

AFFINE AND LINEAR INVARIANT FAMILIES OF HARMONIC MAPPINGS

MARTIN CHUAQUI, RODRIGO HERNÁNDEZ, AND MARÍA J. MARTÍN

ABSTRACT. We study the order of affine and linear invariant families of planar harmonic mappings in the unit disk and determine the order of the family of mappings with bounded Schwarzian norm. The result shows that finding the order of the class \mathcal{S}_H of univalent harmonic mappings can be formulated as a question about Schwarzian norm and, in particular, our result shows consistency between the conjectured order of \mathcal{S}_H and the Schwarzian norm of the harmonic Koebe function.

INTRODUCTION

The purpose of this paper is to study certain affine and linear invariant families of planar harmonic mappings defined in the unit disk \mathbb{D} , with a special interest in the family of mappings with bounded Schwarzian norm. The fundamental aspects of linear invariant families of harmonic mappings were studied in [20], while linear invariant families of holomorphic mappings were introduced by Pommerenke in [17]. Several important properties of such families of either holomorphic or harmonic mappings depend on its order, namely the optimal bound for the second Taylor coefficient of the holomorphic part of the mappings considered. The order of the class of holomorphic mappings with bounded Schwarzian norm can be determined by means of a variational method that gives a relation for the second and third order coefficients of an extremal mapping [17]. Our motivation in this paper stems from the still unresolved problem of determining the order of the family \mathcal{S}_H of normalized univalent harmonic mappings. In this direction, we are able to apply the variational approach that leads to the Marty relations to determine the order of the family of harmonic mappings with a given bound for the Schwarzian derivative. This seems relevant because the order of the class \mathcal{S} of normalized univalent holomorphic mappings can be derived

Date: May 21, 2014.

2010 Mathematics Subject Classification. 30C55.

Key words and phrases. Univalent function, harmonic mapping, Schwarzian derivative, affine and linear invariant family, order.

The authors are partially supported by grants Fondecyt 1110160 and 1110321, Chile. The second author also thankfully acknowledges partial support from Faculty of Forestry and Sciences UEF, Finland (930349). The third author is supported by Academy of Finland grant 268009 and by Spanish MINECO Research Project MTM2012-37436-C02-02.

from the above coefficient relation and the well-known Schwarzian bound for the class. Our result shows consistency between the conjectured values for order in \mathcal{S}_H and in \mathcal{S}_H^0 and the Schwarzian norm of the harmonic Koebe function, a natural candidate for maximizing the Schwarzian norm in \mathcal{S}_H^0 . The hyperbolic norm of the dilatations of the harmonic mappings enter in our analysis in an unexpected way, and turn out to be intimately related to the order of the family. The construction of an extremal mapping for our main result is not elementary and depends on a subtle interplay between the second coefficient and the hyperbolic norm. In another direction, we also establish the sharp bound for the Schwarzian norm of certain important families of harmonic mappings for which the order was already known.

The *Schwarzian derivative* of a locally univalent analytic function on a domain in the complex plane is

$$Sf = \left(\frac{f''}{f'} \right)' - \frac{1}{2} \left(\frac{f''}{f'} \right)^2.$$

The role of Sf in the study of univalence, distortion, and extensions of f has been developed extensively in the literature (see, *e.g.*, [2, 6, 12, 15]). Two of the main properties of the Schwarzian derivative are the following:

- i) $Sf = 0$ if and only if f is a Möbius transformation.
- ii) Whenever the composition $f \circ g$ is well-defined, the *chain rule* holds:

$$S(f \circ g) = (Sf \circ g) \cdot (g')^2 + Sg.$$

In the case when f is locally univalent in the unit disk \mathbb{D} , the *Schwarzian norm*

$$\|Sf\| = \sup_{|z| < 1} |Sf(z)| \cdot (1 - |z|^2)^2$$

turns out to be invariant under post-compositions with automorphisms σ of the disk. In other words, for any such functions f and σ ,

$$\|S(f \circ \sigma)\| = \|Sf\|.$$

In [17, 18] Pommerenke studied and carried through a detailed analysis of the so-called *linear invariant families*; that is, families of locally univalent holomorphic functions f in the unit disk normalized by the conditions $f(0) = 1 - f'(0) = 0$ and which are closed under the transformation

$$F_\zeta(z) = \frac{f\left(\frac{\zeta + z}{1 + \bar{\zeta}z}\right) - f(\zeta)}{(1 - |\zeta|^2)f'(\zeta)}, \quad \zeta \in \mathbb{D}.$$

Several important properties, such as growth, covering, and distortion are determined by the *order* of a linear invariant family \mathcal{F} defined by

$$\alpha(\mathcal{F}) = \sup_{f \in \mathcal{F}} a_2(f) = \frac{1}{2} \sup_{f \in \mathcal{F}} |f''(0)|.$$

For example, the order of the important class \mathcal{S} of normalized univalent mappings in \mathbb{D} is 2. We refer the reader to the books [4] or [19] for more details related to the class \mathcal{S} .

In [17], Pommerenke proves the following theorem regarding the linear invariant family \mathcal{H}_λ of normalized locally univalent analytic functions f in the unit disk with $\|Sf\| \leq \lambda$.

THEOREM A. *The order of the family \mathcal{H}_λ is given by*

$$\alpha(\mathcal{H}_\lambda) = \sqrt{1 + \frac{\lambda}{2}}.$$

It is a straightforward calculation to show that given $\lambda \geq 0$, the function

$$(1) \quad \varphi_a(z) = \frac{1}{2a} \left[\left(\frac{1+z}{1-z} \right)^a - 1 \right], \quad |z| < 1, \quad a = \sqrt{\frac{\lambda}{2} + 1},$$

belongs to \mathcal{H}_λ and satisfies

$$\frac{1}{2}|\varphi_a''(0)| = \sqrt{1 + \frac{\lambda}{2}}.$$

Whenever $f \in \mathcal{S}$, its Schwarzian norm is bounded by 6. Therefore, $\mathcal{S} \subset \mathcal{H}_6$ so that by Theorem A we get $|a_2| \leq 2$ for all $f \in \mathcal{S}$, although the bound will hold also for the non-univalent mappings in \mathcal{H}_6 .

A planar *harmonic mapping* in a domain $\Omega \subset \mathbb{C}$ is a complex-valued function $w = f(z) = u(z) + iv(z)$, $z = x + iy$, which is harmonic, that is, $\Delta u = \Delta v = 0$. When Ω is simply connected, the mapping f has a canonical decomposition $f = h + \bar{g}$, where h and g are analytic in Ω . As is usual, we call h the *analytic part of f* and g the *co-analytic part of f* . The harmonic mapping f is analytic if and only if g is constant. Lewy [13] proved that a harmonic mapping is locally univalent in a domain Ω if and only if its Jacobian does not vanish. In terms of the canonical decomposition $f = h + \bar{g}$, the Jacobian is given by $|h'|^2 - |g'|^2$, and thus, a locally univalent harmonic mapping in a simply connected domain Ω will be *sense-preserving* or *sense-reversing* according to whether $|h'| > |g'|$ or $|g'| > |h'|$ in Ω . Note that $f = h + \bar{g}$ is sense-preserving in Ω if and only if h' does not vanish and the (second complex) *dilatation* $\omega = g'/h'$ has the property that $|\omega| < 1$ in Ω . We refer the reader to the book by Duren [5] for an excellent exposition of harmonic mappings.

Let \mathcal{F} be a family of sense-preserving harmonic mappings $f = h + \bar{g}$ in \mathbb{D} , normalized with $h(0) = g(0) = 0$ and $h'(0) = 1$. The family is said to be *affine and linear invariant* (\mathcal{AL} family) if it closed under the two operations of *Koebe transform* and *affine change*:

$$(2) \quad K_\zeta(f)(z) = \frac{f\left(\frac{z+\zeta}{1+\bar{\zeta}z}\right) - f(\zeta)}{(1-|\zeta|^2)h'(\zeta)}, \quad |\zeta| < 1,$$

and

$$(3) \quad A_\varepsilon(f)(z) = \frac{f(z) - \overline{\varepsilon f(z)}}{1 - \overline{\varepsilon} g'(0)}, \quad |\varepsilon| < 1.$$

Sheil-Small [20] offers an in depth study of affine and linear invariant families \mathcal{F} of harmonic mappings in \mathbb{D} . The *order* of the \mathcal{AL} family, given by

$$\alpha(\mathcal{F}) = \sup_{f \in \mathcal{F}} |a_2(f)| = \frac{1}{2} \sup_{f \in \mathcal{F}} |h''(0)|,$$

plays once more a special role in the analysis.

A special example of affine and linear invariant family is the class \mathcal{S}_H of (normalized) sense-preserving harmonic mappings which are univalent in the unit disk. As it is usual, we use \mathcal{S}_H^0 to denote the family of functions $f = h + \bar{g} \in \mathcal{S}_H$ with $g'(0) = 0$. It is conjectured that the second Taylor coefficient of the analytic part h of any function in \mathcal{S}_H^0 is bounded by $5/2$. If this conjecture were true, we would obtain that the order of \mathcal{S}_H is equal to 3. The analytic part of the so-called harmonic Koebe function $K \in \mathcal{S}_H^0$ (introduced by Clunie and Sheil-Small in [3]) has second coefficient equal to $5/2$.

In [10], the authors introduce a definition for the Schwarzian derivative S_f of locally univalent harmonic mappings, which serves as a complement to the definition found in [1]. The requirement of the latter that the dilatation be a square has been replaced in the former by the local univalence. In both cases a chain rule is in order, which shows that for mappings in \mathbb{D} the norm $\|S_f\|$ defined as before is invariant under automorphisms of the disk. But only the Schwarzian S_f introduced in [10] is invariant under affine changes $af + b\bar{f}$, $|a| \neq |b|$. As a result, the family \mathcal{F}_λ of sense-preserving harmonic mappings $f = h + \bar{g}$ in \mathbb{D} , with $h(0) = g(0) = 0$, $h'(0) = 1$ and $\|S_f\| \leq \lambda$, is affine and linear invariant. We let $\mathcal{F}_\lambda^0 = \{f \in \mathcal{F}_\lambda : g'(0) = 0\}$.

The hyperbolic norm of the dilatation ω of a sense-preserving harmonic mapping in \mathbb{D} is given by

$$\|\omega^*\| = \sup_{z \in \mathbb{D}} \frac{|\omega'(z)| \cdot (1 - |z|^2)}{1 - |\omega(z)|^2},$$

and will play a distinguished role in our analysis. Observe that $\|\omega^*\| \leq 1$ by Schwarz's lemma. As we will show, only for $\lambda \geq 3/2$ will there exist a mapping in $f \in \mathcal{F}_\lambda$ with a dilatation of hyperbolic norm equal to 1. Moreover, for $\lambda < 3/2$ the supremum of all hyperbolic norms of mappings in \mathcal{F}_λ will be strictly less than 1.

Let us denote by \mathcal{A}_λ^0 (resp. \mathcal{A}_λ) the set of *admissible dilatations* of functions $f \in \mathcal{F}_\lambda^0$ (resp. \mathcal{F}_λ); i.e., $\omega \in \mathcal{A}_\lambda^0$ (or \mathcal{A}_λ) if there exists a harmonic mapping $f = h + \bar{g} \in \mathcal{F}_\lambda^0$ (\mathcal{F}_λ) with dilatation ω . The main purpose of this article is to show the following generalization of Theorem A.

THEOREM 1. *The order of \mathcal{F}_λ is given by*

$$\begin{aligned} (4) \quad \alpha(\mathcal{F}_\lambda) &= \sqrt{\frac{\lambda}{2} + 1 + \frac{1}{2} \sup_{f \in \mathcal{F}_\lambda^0} |g''(0)|^2} + \frac{1}{2} \sup_{f \in \mathcal{F}_\lambda^0} |g''(0)| \\ (5) \quad &= \sqrt{\frac{\lambda}{2} + 1 + \frac{1}{2} \sup_{\omega \in \mathcal{A}_\lambda} \|\omega^*\|^2} + \frac{1}{2} \sup_{\omega \in \mathcal{A}_\lambda} \|\omega^*\|. \end{aligned}$$

Furthermore,

$$\begin{aligned} (6) \quad \frac{1}{2} \sup_{f \in \mathcal{F}_\lambda^0} |h''(0)| &= \sqrt{\frac{\lambda}{2} + 1 + \frac{1}{2} \sup_{f \in \mathcal{F}_\lambda^0} |g''(0)|^2} \\ (7) \quad &= \sqrt{\frac{\lambda}{2} + 1 + \frac{1}{2} \sup_{\omega \in \mathcal{A}_\lambda^0} |\omega'(0)|^2}. \end{aligned}$$

It was proved in [10] that the Schwarzian norm of the harmonic Koebe function equals $19/2$. For $\lambda = 19/2$ our result gives

$$\frac{1}{2} \sup_{f \in \mathcal{F}_\lambda^0} |h''(0)| = \frac{5}{2},$$

which would show that the order of \mathcal{S}_H is equal to 3 provided the harmonic Koebe function was extremal in the class for the Schwarzian norm.

1. SCHWARZIAN DERIVATIVE

Let $f = h + \bar{g}$ be a locally univalent harmonic mapping in a simply connected domain Ω with dilatation $\omega = g'/h'$. In [10], the Schwarzian derivative S_f of such a function f was defined. If f is sense-preserving, S_f is given by

$$(8) \quad S_f = Sh + \frac{\bar{\omega}}{1 - |\omega|^2} \left(\frac{h''}{h'} \omega' - \omega'' \right) - \frac{3}{2} \left(\frac{\omega' \bar{\omega}}{1 - |\omega|^2} \right)^2.$$

Several properties of this operator are the following:

- (i) $S_f \equiv 0$ if and only if $f = \alpha T + \beta \bar{T}$, where $|\alpha| \neq |\beta|$ and T is a Möbius transformation of the form

$$T(z) = \frac{az + b}{cz + d}, \quad ad - bc \neq 0.$$

- (ii) Whenever f is a sense-preserving harmonic mapping and ϕ is an analytic function such that the composition $f \circ \phi$ is well-defined, the Schwarzian derivative of $f \circ \phi$ can be computed using the *chain rule*

$$S_{f \circ \phi} = S_f(\phi) \cdot (\phi')^2 + S\phi.$$

- (iii) For any affine mapping $L(z) = az + b\bar{z}$ with $|a| \neq |b|$, we have that $S_{L \circ f} = S_f$. Note that L is sense-preserving if and only if $|b| < |a|$.

Consider now a sense-preserving harmonic mapping f in the unit disk. Using the chain rule, we can see that for each $z \in \mathbb{D}$

$$|S_f(z)| = |S_{(f \circ \sigma_z)}(0)| \cdot (1 - |z|^2)^2,$$

where σ_z is any automorphism of the unit disk with $\sigma_z(0) = z$. The *Schwarzian norm* $\|S_f\|$ of f is defined by

$$\|S_f\| = \sup_{z \in \mathbb{D}} |S_f(z)| \cdot (1 - |z|^2)^2.$$

It is easy to check (using the chain rule again and the Schwarz-Pick lemma) that $\|S_{f \circ \sigma}\| = \|S_f\|$ for any automorphism of the unit disk σ . For further properties of S_f and the motivation for this definition, see [10].

2. AFFINE AND LINEAR INVARIANT FAMILIES

Let \mathcal{S}_H denote the family of sense-preserving univalent harmonic mappings $f = h + \bar{g}$ on \mathbb{D} normalized by $h(0) = 0$, $h'(0) = 1$, and $g(0) = 0$. This family is affine and linear invariant. As usual, we use \mathcal{S}_H^0 to denote the subclass of functions in \mathcal{S}_H that satisfy the further normalization $g'(0) = 0$. The family \mathcal{S}_H is normal and \mathcal{S}_H^0 is compact (see [3] or [5]). Analogous results are obtained when dealing with the families \mathcal{C}_H and \mathcal{C}_H^0 of convex harmonic mappings in \mathcal{S}_H and \mathcal{S}_H^0 , respectively.

Other examples of \mathcal{AL} families of sense-preserving harmonic mappings are the stable harmonic univalent (\mathcal{SHU}) and the stable harmonic convex (\mathcal{SHC}) classes. A function $f = h + \bar{g} \in \mathcal{S}_H$ is \mathcal{SHU} (resp. \mathcal{SHC}) if $h + \lambda \bar{g}$ is univalent (convex) for every $|\lambda| = 1$. These classes are linear invariant, and also affine because univalence or convexity are preserved under the affine changes A_ε as in (3). An important observation is that if the harmonic mapping f has dilatation ω , then $F = A_{\omega(0)}(f)$ will have a dilatation vanishing at the origin.

It is easy to check that $\alpha(\mathcal{SHU}) = 2$ and $\alpha(\mathcal{SHC}) = 1$ (see [8]). In the next theorem, we obtain sharp bounds for the Schwarzian norm of functions in these classes.

THEOREM 2. *Let $f = h + \bar{g}$ be a locally univalent harmonic mapping defined in \mathbb{D} .*

- (i) *If f is a \mathcal{SHU} mapping, then $\|S_f\| \leq 6$.*
- (ii) *If f is a \mathcal{SHC} mapping, then $\|S_f\| \leq 2$.*

Both constants are sharp.

Proof. It was shown in [8] that if $f = h + \bar{g} \in \mathcal{SHU}$, then $h + ag$ is univalent for all $|a| < 1$; in particular, h itself is univalent. Assume that there exists $f \in \mathcal{SHU}$ with $\|S_f\| > 6$ and let ω be its dilatation. Then, there is a point $\zeta \in \mathbb{D}$ such that $|S_f(\zeta)| \cdot (1 - |\zeta|^2)^2 > 6$. Using the chain rule for the Schwarzian derivative and the affine invariance, we see that

$$(9) \quad |S_{A_{\omega(\zeta)}(K_\zeta(f))}(0)| = |S_f(\zeta)| \cdot (1 - |\zeta|^2)^2 > 6,$$

where K_ζ is the transformation defined by (2), ω_ζ is the dilatation of the function $K_\zeta(f)$, and $A_{\omega_\zeta(0)}$ is as in (3). Since the dilatation of $K_\zeta(f)$ at the origin is $\omega_\zeta(0)$, we have that the dilatation of $A_{\omega_\zeta(0)}(K_\zeta(f))$ fixes the origin. Let H denote the (univalent) analytic part of $A_{\omega_\zeta(0)}(K_\zeta(f))$; keeping in mind the definition (8) for the Schwarzian derivative, we see from (9) that

$$|S_{A_{\omega_\zeta(0)}(K_\zeta(f))}(0)| = |SH(0)| > 6,$$

which contradicts the univalence of H . This proves statement (i). The proof of (ii) follows the same argument, except for the fact that convex analytic mappings have Schwarzian norm bounded by 2 [16].

To prove that both constants are sharp, it is enough to consider the analytic functions

$$k(z) = \frac{z}{(1-z)^2} \quad \text{and} \quad s(z) = \frac{1}{2} \log \left(\frac{1+z}{1-z} \right)$$

that belong to the families of \mathcal{SHU} and \mathcal{SHC} mappings and have Schwarzian norms $\|S_k\| = 6$ and $\|S_s\| = 2$, respectively. \square

3. ADMISSIBLE DILATATIONS

In this section, we review some of the properties of hyperbolic derivatives of self-maps of the unit disk and determine the relation between hyperbolic norms of admissible dilatations in \mathcal{F}_λ and the parameter λ itself.

3.1. The hyperbolic derivative. Let ω be a self-map of the unit disk, this is, an analytic function in \mathbb{D} with $\omega(\mathbb{D}) \subset \mathbb{D}$. The *hyperbolic derivative* of such function ω is

$$\omega^*(z) = \frac{\omega'(z) \cdot (1 - |z|^2)}{1 - |\omega(z)|^2}, \quad z \in \mathbb{D}.$$

From Schwarz's lemma we see that $|\omega^*| \leq 1$ in \mathbb{D} , and that if there exists $z_0 \in \mathbb{D}$ with $|\omega^*(z_0)| = 1$, then ω is an automorphism of the unit disk and $|\omega^*| \equiv 1$ in \mathbb{D} . Of course, there are self-maps of \mathbb{D} with hyperbolic norm equal to 1 which are not automorphisms. There are examples such as $\omega(z) = (z+1)/2$, but also, every finite Blaschke product has hyperbolic norm equal to 1 [7]. See also [14] for other examples.

Given two self-maps ω and φ of the unit disk, the *chain rule for the hyperbolic derivative* holds:

$$(\varphi \circ \omega)^*(z) = \varphi^*(\omega(z)) \cdot \omega^*(z).$$

In particular, if σ is an automorphism of \mathbb{D} , then $|(\sigma \circ \omega)^*| \equiv |\omega^*|$ in the unit disk, hence $\|(\sigma \circ \omega)^*\| = \|\omega\|$.

Just like in [9], the so-called *lens-maps* ℓ_α will be of particular interest. For $0 < \alpha < 1$, the mapping ℓ_α is defined by

$$(10) \quad \ell_\alpha(z) = \frac{\ell(z)^\alpha - 1}{\ell(z)^\alpha + 1},$$

where $\ell(z) = (1+z)/(1-z)$. The hyperbolic norm of ℓ_α was computed in [11] to be $\|\ell_\alpha^*\| = \alpha$. Moreover, $|\ell_\alpha^*(r)| = \alpha$ for all real numbers $0 \leq r < 1$.

3.2. Admissible dilatations and norms. Recall that we say that a self-map ω of the unit disk belongs to the family of admissible dilatations \mathcal{A}_λ^0 if there exists $f \in \mathcal{F}_\lambda^0$ with dilatation ω . For any such f and any $\alpha \in \mathbb{D}$, the affine transformation $f_\alpha = A_{-\alpha}(f) = f + \overline{\alpha}f \in \mathcal{F}_\lambda$. Its dilatation ω_α is given by

$$\omega_\alpha = \sigma_\alpha \circ \omega,$$

where σ_α is the automorphism in the unit disk defined by

$$(11) \quad \sigma_\alpha(z) = \frac{\alpha + z}{1 + \overline{\alpha}z}, \quad z \in \mathbb{D}.$$

In other words, whenever $\omega \in \mathcal{A}_\lambda^0$ and $\alpha \in \mathbb{D}$, then $\sigma_\alpha \circ \omega \in \mathcal{A}_\lambda$.

On the other hand, if $f \in \mathcal{F}_\lambda$ that has dilatation ω satisfying $\omega(0) = \alpha$, then

$$F = \frac{f - \overline{\alpha}f}{1 - |\alpha|^2} \in \mathcal{F}_\lambda^0,$$

with a resulting new dilatation $\omega_F = \sigma_\alpha^{-1} \circ \omega = \sigma_{-\alpha} \circ \omega$. These relations establish a correspondence between the families \mathcal{A}_λ^0 and \mathcal{A}_λ .

It is easy to verify that for any value of λ , there exists $\omega \in \mathcal{A}_\lambda$ with $\|\omega^*\| = 0$ (just consider the identity function $I(z) = z$ in the unit disk which belongs to \mathcal{F}_λ for all $\lambda \geq 0$). The following theorem characterizes the values of λ for which a dilatation in \mathcal{A}_λ can have hyperbolic norm 1.

THEOREM 3. *The following conditions are equivalent.*

- (i) $\lambda \geq 3/2$.
- (ii) *There exists $\omega \in \mathcal{A}_\lambda^0$ with $|\omega'(0)| = 1$.*
- (iii) *The set $\{\lambda \cdot I : |\lambda| = 1\}$ is contained in \mathcal{A}_λ^0 . In particular, the identity function I is an admissible dilatation in \mathcal{F}_λ^0 .*
- (iv) *Every automorphism σ of the unit disk is an admissible dilatation in \mathcal{F}_λ .*
- (v) *There exist $\omega \in \mathcal{A}_\lambda$ with $\|\omega^*\| = 1$.*

Proof. The scheme of the proof is to show that (i) \Rightarrow (ii) \Rightarrow (iii) \Rightarrow (i). Then we prove that (iii) \Longleftrightarrow (iv), and finally we see that (iv) \Rightarrow (v) \Rightarrow (ii).

To show that (i) \Rightarrow (ii), we consider the function $f = z + \frac{1}{2}\overline{z}^2$. Note that the dilatation ω of f equals the identity function I so that $\omega'(0) = 1$. Since

$$S_f(z) = -\frac{3\overline{z}^2}{2(1-|z|^2)^2},$$

we have $\|S_f\| = 3/2$. Thus, $f \in \mathcal{F}_\lambda^0$ for all $\lambda \geq 3/2$ and (ii) holds.

Let us now check that (ii) \Rightarrow (iii). To do so, just note that given any harmonic function $f = h + \overline{g} \in \mathcal{F}_\lambda^0$ with dilatation ω and any $\mu \in \partial\mathbb{D}$, the functions $f_\mu = h + \overline{\mu g}$ belong to \mathcal{F}_λ^0 as well since $S_f = S_{f_\mu}$ for all such μ and the dilatation ω_μ of f_μ is equal to $\omega_\mu(z) = \mu \omega(z)$ (which satisfies $\omega_\mu(0) = 0$).

Now, if we assume that there exists $\omega \in \mathcal{A}_\lambda^0$ with $|\omega'(0)| = 1$ we have, by the Schwarz lemma, that ω is a rotation of the disk (i.e. $\omega(z) = \lambda z$ for some $|\lambda| = 1$). Therefore, using those functions f_μ we immediately get that (iii) holds.

Suppose that (iii) is satisfied. This is, there is a function $f = h + \bar{g} \in \mathcal{F}_\lambda^0$ with dilatation $\omega = I$. We are going to prove that $\lambda \geq 3/2$. Indeed, we will obtain a stronger result in view of inequality (16) below. More specifically, we will see that (iii) implies that the order $\alpha(\mathcal{F}_\lambda)$ of the family \mathcal{F}_λ is 2 at least, so that by (16) we obtain (i) (note that the supremum that appears there equals 1 whenever $I \in \mathcal{A}_\lambda^0$).

Using the affine invariance property of \mathcal{F}_λ , we have that whenever $f \in \mathcal{F}_\lambda^0$, the function $F_\varepsilon = A_\varepsilon(f) = f - \varepsilon \bar{f} \in \mathcal{F}_\lambda$ for all $\varepsilon \in \mathbb{D}$, where A_ε is the transformation defined in (3). The analytic part in the canonical decomposition of F_ε equals $h_\varepsilon = h - \varepsilon \bar{g}$ so that $h'_\varepsilon(z) = h'(z) \cdot (1 - \varepsilon \bar{I})$. Therefore, just applying the same arguments that appear in the proof of the theorem on [5, p. 97], we get that for all $z \in \mathbb{D}$

$$(12) \quad |h'(z)| \cdot |1 - \varepsilon z| \geq \frac{(1 - |z|)^{\alpha(\mathcal{F}_\lambda)-1}}{(1 + |z|)^{\alpha(\mathcal{F}_\lambda)+1}}$$

whenever $\varepsilon < 1$ and hence, for all $|\varepsilon| \leq 1$. Thus given any $z \neq 0$ in the unit disk, we can choose $\varepsilon = \bar{z}/|z|$ to get from (12) that

$$|h'(z)| \geq \frac{(1 - |z|)^{\alpha(\mathcal{F}_\lambda)-2}}{(1 + |z|)^{\alpha(\mathcal{F}_\lambda)+1}},$$

an inequality that obviously works for $z = 0$. As a consequence, since h is locally univalent in \mathbb{D} , we see that the analytic function $1/h'$ in the unit disk satisfies

$$\frac{1}{|h'(z)|} \leq \frac{(1 + |z|)^{\alpha(\mathcal{F}_\lambda)+1}}{(1 - |z|)^{\alpha(\mathcal{F}_\lambda)-2}},$$

which implies, by the maximum modulus principle, that $\alpha(\mathcal{F}_\lambda) \geq 2$ (otherwise we would get $1/|h'(0)| < 1$, which is absurd since $h'(0) = 1$). This shows that (iii) \Rightarrow (i).

We continue with the proof of the theorem by showing that (iii) \Leftrightarrow (iv). Note that since every rotation is an automorphism of the unit disk that fixes the origin, only the implication (iii) \Rightarrow (iv) is needed to check the equivalence between these two statements. Let us assume then that the set $\{\lambda I : |\lambda| = 1\} \subset \mathcal{A}_\lambda^0$. Then, we can use that \mathcal{F}_λ is affine invariant and the functions f_μ defined above to see that for any α in the unit disk and any $\lambda, \mu \in \partial\mathbb{D}$, the functions

$$\mu \sigma_\alpha \circ (\lambda I) = \mu \cdot \frac{\alpha + \lambda z}{1 + \bar{\alpha} \lambda z} = \mu \lambda \cdot \frac{\bar{\lambda} \alpha + z}{1 + \bar{\bar{\lambda}} \alpha z} \in \mathcal{A}_\lambda.$$

In other words, for any $|\eta| = 1$ and any $\beta \in \mathbb{D}$, the mappings

$$(13) \quad \eta \cdot \frac{\beta + z}{1 + \overline{\beta}z} \in \mathcal{A}_\lambda.$$

The Schwarz lemma says that any automorphism of the unit disk has the form (13) so that (iv) holds.

Finally, we show the equivalence between statements (iv) and (v). Once more, only one of the implications in the equivalence is non-trivial (recall that any automorphism of the unit disk has hyperbolic norm equal to one). Concretely, we just need to see that (v) \Rightarrow (iv). Since at this point we have proved that (i), (ii), (iii), and (iv) are equivalent, it suffices to check that (v) \Rightarrow (ii). To do so, let $\omega \in \mathcal{A}_\lambda$ have the property that $\|\omega^*\| = 1$. Then, by the definition of the hyperbolic norm, there exists a sequence of points in the unit disk $\{z_n\}$, say, such that

$$\lim_{n \rightarrow \infty} |\omega^*(z_n)| = 1.$$

The fact that $\omega \in \mathcal{A}_\lambda$ means that there is a function $f = h + \overline{g} \in \mathcal{F}_\lambda$ with dilatation ω . Using the transformations K_{z_n} and $A_{\omega_n(0)}$ defined by (2) and (3), respectively, where ω_n is the dilatation of $K_{z_n}(f)$, we obtain a sequence of harmonic mappings

$$f_n = A_{\omega_n(0)}(K_{z_n}(f))$$

with dilatations

$$\gamma_n = \sigma_{-\omega_n(0)} \circ \omega_n = \sigma_{-\omega_n(0)} \circ [\lambda_n \cdot (\omega \circ \sigma_{z_n})],$$

where $\lambda_n = h'(z_n)/\overline{h'(z_n)}$ and σ_α is again the automorphism of \mathbb{D} defined by (11).

Note that $\gamma_n(0) = 0$, so that each element in the sequence $f_n \in \mathcal{F}_\lambda^0$. Moreover, a straightforward computation shows that we also have $|\gamma'_n(0)| = |\omega^*(z_n)|$.

Now, we argue as on [5, pp. 81-82] but in this case, instead of the argument principle, we use the fact that whenever $f_n \rightarrow f_0$ uniformly on compact subsets of \mathbb{D} we have that for each z in the unit disk

$$S_{f_n}(z) \rightarrow S_{f_0}(z).$$

This shows that \mathcal{F}_λ^0 is a normal and compact family. Thus, there exists a subsequence (that we rename $\{f_n\}$ again) of the sequence $\{f_n\}$ that converges to $f_0 \in \mathcal{F}_\lambda^0$ uniformly on compact subsets in the unit disk. The dilatation γ_0 of the limit function f_0 is

$$\gamma_0(z) = \lim_{n \rightarrow \infty} \gamma_n(z),$$

so that it satisfies

$$|\gamma'_0(0)| = \left| \lim_{n \rightarrow \infty} \gamma'_n(0) \right| = \lim_{n \rightarrow \infty} |\omega^*(z_n)| = 1.$$

This proves that (v) \Rightarrow (ii) and finishes the proof of the theorem. \square

From Theorems 1 and 3 we derive the following important corollary.

COROLLARY 1. *If $\lambda \geq 3/2$ then*

$$\alpha(\mathcal{F}_\lambda) = \sqrt{\frac{\lambda}{2} + \frac{3}{2}} + \frac{1}{2}$$

and

$$\frac{1}{2} \sup_{f \in \mathcal{F}_\lambda^0} |h''(0)| = \sqrt{\frac{\lambda}{2} + \frac{3}{2}}.$$

Even though for $\lambda < 3/2$ we have been unable to determine the value of $\sup_{\omega \in \mathcal{A}_\lambda} \|\omega^*\|$, one can rewrite Theorem 1 in the form

$$(14) \quad \sup_{\omega \in \mathcal{A}_\lambda} \|\omega^*\| = -2\alpha(\mathcal{F}_\lambda) + \sqrt{8\alpha(\mathcal{F}_\lambda)^2 - 2\lambda - 4}.$$

4. PROOF OF THE MAIN THEOREM

We will divide the proof of Theorem 1 in three different parts.

4.1. **A lemma.** We begin by proving that (5) and (7) hold.

LEMMA 1. *For any positive real number λ ,*

(i)

$$\sup_{f=h+\bar{g} \in \mathcal{F}_\lambda^0} |g''(0)| = \sup_{\omega \in \mathcal{A}_\lambda^0} |\omega'(0)|.$$

(ii)

$$\sup_{\omega \in \mathcal{A}_\lambda^0} |\omega'(0)| = \sup_{\omega \in \mathcal{A}_\lambda^0} \|\omega^*\| = \sup_{\omega \in \mathcal{A}_\lambda} \|\omega^*\|.$$

Proof. The proof of (i) is trivial: just recall that the dilatation ω of $f = h + \bar{g} \in \mathcal{F}_\lambda^0$ equals $\omega = g'/h'$ and satisfies $\omega(0) = 0$. Thus, $g' = \omega h'$ and hence $g''(0) = \omega'(0)h'(0) + \omega(0)h''(0) = \omega'(0)$, since $h'(0) = 1$.

To prove (ii), take an arbitrary $f = h + \bar{g} \in \mathcal{F}_\lambda^0$ with dilatation ω . The transformations (2) and (3) produce new functions that also belong to this family. Concretely, given any such function f , we define the mappings

$$f_\zeta = A_{\omega_\zeta(0)}(K_\zeta(f)) = h_\zeta + \bar{g}_\zeta, \quad \zeta \in \mathbb{D},$$

where ω_ζ is the dilatation of $K_\zeta(f)$. These mappings f_ζ belong to \mathcal{F}_λ^0 and have dilatations

$$\gamma_\zeta = \sigma_{-\omega_\zeta(0)} \circ \omega_\zeta \in \mathcal{A}_\lambda^0,$$

where $\omega_\zeta = \lambda_\zeta \cdot (\omega \circ \sigma_\zeta)$ for certain $|\lambda_\zeta| = 1$ and appropriate automorphisms σ_α of the form (11). Since

$$|\gamma'_\zeta(0)| = |\omega^*(\zeta)|$$

and ζ is any arbitrary point in \mathbb{D} , we conclude that the first equality in (ii) holds. The second inequality is an easy consequence of the fact that the correspondence between the families \mathcal{A}_λ and \mathcal{A}_λ^0 is realized using precomposition with automorphisms of the unit disk (an operation that preserves the hyperbolic norm). \square

4.2. Upper bounds. The aim of this section is to prove the following inequalities:

$$(15) \quad \frac{1}{2} \sup_{f \in \mathcal{F}_\lambda^0} |h''(0)| \leq \sqrt{\frac{\lambda}{2} + 1} + \frac{1}{2} \sup_{\omega \in \mathcal{A}_\lambda^0} |\omega'(0)|^2$$

and

$$(16) \quad \alpha(\mathcal{F}_\lambda) \leq \sqrt{\frac{\lambda}{2} + 1} + \frac{1}{2} \sup_{\omega \in \mathcal{A}_\lambda} \|\omega^*\|^2 + \frac{1}{2} \sup_{\omega \in \mathcal{A}_\lambda} \|\omega^*\|.$$

To prove (15), let us agree with the notation

$$S = \frac{1}{2} \sup_{f \in \mathcal{F}_\lambda^0} |h''(0)|.$$

Given $f \in \mathcal{F}_\lambda^0$, the function F defined by $F(z) = \bar{\lambda}f(\lambda z) \in \mathcal{F}_\lambda^0$ as well. Therefore, we see that

$$S = \frac{1}{2} \sup_{f \in \mathcal{F}_\lambda^0} \operatorname{Re} \{h''(0)\}.$$

As it was mentioned in the proof of Theorem 3, the family \mathcal{F}_λ^0 is normal and compact. Therefore, there exists a function $f_0 = h_0 + \overline{g_0}$ in \mathcal{F}_λ^0 with dilatation ω_0 , where

$$h_0(z) = z + a_2 z^2 + a_3 z^3 + \dots, \quad g_0(z) = b_2 z^2 + b_3 z^3 + \dots,$$

such that

$$\operatorname{Re} \{a_2\} = |a_2| = S.$$

Take an arbitrary point $\zeta \in \mathbb{D}$ and consider, once more, the transformations (2) and (3) to produce the family

$$F_\zeta = A_{\omega_\zeta(0)}(K_\zeta(f)) = h_\zeta^* + \overline{g_\zeta^*}$$

of functions that are in \mathcal{F}_λ^0 .

As it is shown on [5, p.102], the Taylor coefficients a_n^* of h_ζ^* satisfy

$$a_n^* = a_n + [(n+1)a_{n+1} - 2a_2 a_n] \zeta - [2\overline{b_2} b_n + (n-1)a_{n-1}] \overline{\zeta} + o(|\zeta|).$$

Then, we get that

$$(17) \quad 3a_3 - 2a_2^2 - 2|b_2|^2 - 1 = 0.$$

Now, using the equations $S_{f_0}(0) = S h_0(0) = 6(a_3 - a_2^2)$, we have by (17) that $\operatorname{Re} \{S_{f_0}(0)\} = 4|b_2|^2 + 2 - 2\operatorname{Re} \{a_2^2\}$. Hence, bearing in mind $S = \operatorname{Re} \{a_2\} = |a_2|$, we obtain

$$\begin{aligned} S^2 &= |a_2|^2 \leq \frac{1}{2} |S_{f_0}(0)| + 1 + 2|b_2|^2 \\ &\leq \frac{\lambda}{2} + 1 + \frac{1}{2} \sup_{\omega \in \mathcal{A}_\lambda^0} |\omega'(0)|^2. \end{aligned}$$

This proves (15). To show that (16) holds we can argue as follows. Take any function $f = h + \overline{g} \in \mathcal{F}_\lambda$ with dilatation ω and consider the affine transformation $A_{\omega(0)}$ as in (3) to get the function

$$F = A_{\omega(0)}(f) = H + \overline{G} \in \mathcal{F}_\lambda^0.$$

Note that $f = F + \overline{\omega(0)F}$. Thus, the analytic part of f equals $H + \overline{\omega(0)G}$. Using also that $\omega(\mathbb{D}) \subset \mathbb{D}$, and Lemma 1, we get

$$\begin{aligned} |h''(0)| &= |H''(0) + \overline{\omega(0)G''(0)}| \leq 2S + |\omega(0)| \sup_{f \in \mathcal{F}_\lambda^0} |g''(0)| \\ &\leq 2S + \sup_{\omega \in \mathcal{A}_\lambda} \|\omega^*\|, \end{aligned}$$

which gives (16).

4.3. Equalities. In order finish the proof of Theorem 1, we will show that (15) and (16) are actual equalities.

To see that equality holds in (15), we will exhibit a function $f_0 = h_0 + \overline{g_0} \in \mathcal{F}_\lambda^0$ with dilatation ω_0 satisfying the following properties:

$$\omega_0'(0) = \sup_{\omega \in \mathcal{A}_\lambda^0} \|\omega^*\| \quad \text{and} \quad h_0''(0) = 2 \cdot \sqrt{\frac{\lambda}{2} + 1 + \omega_0'(0)}.$$

This will show that (6) holds. Furthermore, for any $0 < r < 1$, the mappings $f_r = f_0 + r\overline{f_0} \in \mathcal{F}_\lambda$ and have analytic parts $h_r = h_0 + rg_0$, for which

$$\begin{aligned} h_r''(0) &= h_0''(0) + rg_0''(0) = \sup_{f \in \mathcal{F}_\lambda^0} |h''(0)| + r\omega_0'(0) \\ &= \sup_{f \in \mathcal{F}_\lambda^0} |h''(0)| + r \sup_{\omega \in \mathcal{A}_\lambda^0} \|\omega^*\| \\ &\rightarrow \sup_{f \in \mathcal{F}_\lambda^0} |h''(0)| + \sup_{\omega \in \mathcal{A}_\lambda^0} \|\omega^*\|, \end{aligned}$$

as $r \rightarrow 1^-$. This shows that equality also holds in (16), proving thus (4).

The function f_0 will be constructed in the next, final section.

5. THE MAPPING f_0

Given $\lambda \geq 0$, let us denote by $R_\lambda = \sup_{\omega \in \mathcal{A}_\lambda} \|\omega^*\|$. If $\lambda = 0$, then by [10, Cor. 2] every function in \mathcal{F}_λ has the form $h + a\overline{h}$, where $a \in \mathbb{D}$ and

$$h(z) = \frac{z}{1+bz}, \quad |b| < 1.$$

Hence $R_\lambda = 0$. We now analyze the cases $\lambda > 0$.

5.1. Harmonic mappings in \mathcal{F}_λ^0 when $\lambda > 0$. By Theorem 3, we have that $\lambda \geq 3/2$ if and only if $R_\lambda = 1$. Equation (14) gives the exact value of R_λ in terms of λ and $\alpha(\mathcal{F}_\lambda)$. The following example allows us to estimate the value of R_λ in terms of λ for $0 < \lambda < 3/2$.

Given $0 < \lambda < 3/2$, we can write $\lambda = 3s^2/2$ for some $0 < s < 1$. Consider the functions $f_r = z + \frac{1}{2}r\bar{z}^2$, where $r \in (0, 1)$. The Schwarzian norm of f_r equals

$$(18) \quad \|S_{f_r}\| = \frac{3}{2} \sup_{z \in \mathbb{D}} \frac{r^2|z|^2}{(1 - |rz|^2)^2} \cdot (1 - |z|^2)^2 \leq \frac{3}{2} r^2.$$

Hence, $f_r \in \mathcal{F}_\lambda^0$ whenever $r \leq s$. This implies that $R_\lambda \geq r = \|\omega_r^*\|$, where ω_r is the dilatation of f_r . In particular, this implies that if $\lambda \in (0, 3/2)$,

$$(19) \quad R_\lambda > \sqrt{\frac{2\lambda}{3}}.$$

Note that we have strict inequality in (19) since the supremum that appears in (18) is strictly less than r^2 . Also, that $R_\lambda > 0$ if $\lambda > 0$.

5.2. The extremal. Let us introduce the notation $\alpha = \alpha(\mathcal{F}_\lambda)$ and $R = R_\lambda = \sup\{\|\omega^*\| : \omega \in \mathcal{A}_\lambda\}$. For $\lambda > 0$ we consider the mapping $f_0 = h_0 + \bar{g}_0$, where h_0 and g_0 solve the linear system of equations

$$(20) \quad \begin{cases} h_0 - g_0 = \varphi_a \\ \omega_0 = g'_0/h'_0 = \ell_R \end{cases}, \quad h_0(0) = g_0(0) = 0,$$

where

$$a = \sqrt{\frac{\lambda}{2} + 1 + \frac{1}{2}R^2} - \frac{R}{2},$$

and the function φ_a is the generalized Koebe function defined by (1). Also, ℓ_1 is the identity mapping I and for $0 < R < 1$, ℓ_R is the lens-map (10). Note that $h'_0(0) = 1 - g'_0(0) = 0$ and that f_0 is a locally univalent mapping in the unit disk since the dilatation is a self-map of \mathbb{D} and h_0 is locally univalent.

By (20), we have $h'_0(1 - \ell_R) = \varphi'_a$ which implies

$$\frac{h''_0}{h'_0} = \frac{\varphi''_a}{\varphi'_a} + \frac{\ell'_R}{1 - \ell_R},$$

so that

$$Sh_0 = S\varphi_a + \frac{\ell''_R}{1 - \ell_R} + \frac{1}{2} \frac{(\ell'_R)^2}{(1 - \ell_R)^2} - \frac{\ell'_R}{1 - \ell_R} \cdot \frac{\varphi''_a}{\varphi'_a},$$

which gives

$$Sh_0(0) = 2(1 - a^2) + \frac{R^2}{2} - 2aR = 2 + \frac{R^2}{2} - 2aR - 2a^2.$$

Define the function

$$\psi(x) = 2 + \frac{R^2}{2} - 2xR - 2x^2.$$

Note that

$$\psi \left(\sqrt{1 + \frac{R^2}{2}} - \frac{R}{2} \right) = 0.$$

Also, that $\psi' < 0$ which implies, since

$$a > \sqrt{1 + \frac{R^2}{2}} - \frac{R}{2},$$

that $\psi(a) < 0$. This means that $|Sh_0(0)| = 2a^2 + 2aR - 2 - R^2/2$, which is easily seen to be equal to λ . Since $\omega_0(0) = 0$, we also have that $|Sf_0(0)| = \lambda$. To show that $f_0 \in \mathcal{F}_\lambda^0$, we just need to check $\lambda = |Sf_0(0)| = \|Sf_0\|$. To do so, we compute the Schwarzian derivative of f_0 which, according to (8), equals

$$\begin{aligned} S_{f_0} &= Sh_0 + \frac{\overline{\ell_R}}{1 - |\ell_R|^2} \left(\frac{h_0''}{h_0'} \ell_R' - \ell_R'' \right) - \frac{3}{2} \left(\frac{\overline{\ell_R} \ell_R'}{1 - |\ell_R|^2} \right)^2 \\ (21) \quad &= S\varphi_a + \frac{\ell_R''}{1 - \ell_R} + \frac{1}{2} \frac{(\ell_R')^2}{(1 - \ell_R)^2} - \frac{\ell_R'}{1 - \ell_R} \cdot \frac{\varphi_a''}{\varphi_a'} \\ &+ \frac{\overline{\ell_R}}{1 - |\ell_R|^2} \left(\frac{\varphi_a''}{\varphi_a'} \ell_R' + \frac{(\ell_R')^2}{1 - \ell_R} - \ell_R'' \right) - \frac{3}{2} \left(\frac{\overline{\ell_R} \ell_R'}{1 - |\ell_R|^2} \right)^2. \end{aligned}$$

Since for any complex number z ,

$$\frac{1}{1 - z} - \frac{\overline{z}}{1 - |z|^2} = \frac{1 - \overline{z}}{(1 - z)(1 - |z|^2)},$$

we get from (21)

$$\begin{aligned} S_{f_0} &= S\varphi_a + \frac{1 - \overline{\ell_R}}{1 - \ell_R} \cdot \frac{\ell_R''}{1 - |\ell_R|^2} - \frac{1 - \overline{\ell_R}}{1 - \ell_R} \cdot \frac{\ell_R'}{1 - |\ell_R|^2} \cdot \frac{\varphi_a''}{\varphi_a'} \\ &+ \frac{(\ell_R')^2}{2} \cdot \left[\left(\frac{1 + \overline{\ell_R} - 2|\ell_R|^2}{(1 - \ell_R)(1 - |\ell_R|^2)} \right)^2 - 4 \cdot \left(\frac{\overline{\ell_R}}{1 - |\ell_R|^2} \right)^2 \right] \\ (22) \quad &= S\varphi_a + F_R \cdot \left(\frac{\ell_R''}{\ell_R'} - \frac{\varphi_a''}{\varphi_a'} \right) + \frac{F_R^2}{2} \cdot \frac{1 + 3\overline{\ell_R} - 4|\ell_R|^2}{1 - \overline{\ell_R}}, \end{aligned}$$

where

$$F_R = \frac{1 - \overline{\ell_R}}{1 - \ell_R} \cdot \frac{\ell_R'}{1 - |\ell_R|^2}.$$

Let us write

$$(23) \quad w = \left(\frac{1 + z}{1 - z} \right)^R \quad \text{and} \quad \beta = \frac{w}{Re w}.$$

Then, $F_R(z) = R\beta/(1 - z^2)$,

$$\frac{\ell_R''(z)}{\ell_R'(z)} - \frac{\varphi_a''(z)}{\varphi_a'(z)} = \frac{-2a}{1 - z^2} + \frac{2R}{1 - z^2} \cdot \frac{1 - w}{1 + w},$$

and

$$\frac{1 + 3\overline{\ell_R} - 4|\ell_R|^2}{1 - \overline{\ell_R}} = \frac{5\operatorname{Re}\{w\} - 3(1 + i\operatorname{Im}\{w\})}{1 + w} = \frac{8\operatorname{Re}\{w\}}{1 + w} - 3.$$

Thus, we get from (22)

$$\begin{aligned} S_{f_0}(z) \cdot (1 - z^2)^2 &= 2(1 - a^2) - 2\alpha R\beta - \frac{3R^2\beta^2}{2} \\ &+ 2R^2\beta \cdot \frac{1 - w}{1 + w} + 4R^2\beta^2 \cdot \frac{\operatorname{Re}\{w\}}{1 + w} \\ &= 2(1 - a^2) + 2R(R - a)\beta - \frac{3R^2\beta^2}{2}, \end{aligned}$$

a formula that we use to show

$$\begin{aligned} (24) \quad &|S_{f_0}(z)| \cdot (1 - |z|^2)^2 \\ &= \left| 2(1 - a^2) + 2R(R - a)\beta - \frac{3R^2\beta^2}{2} \right| \cdot \left(\frac{1 - |z|^2}{|1 - z^2|} \right)^2. \end{aligned}$$

Consider a real number γ with $0 \leq \gamma < \pi/2$ and define the curves

$$C_\gamma = \left\{ z \in \mathbb{D} : \operatorname{Arg} \left(\frac{1 + z}{1 - z} \right) = \gamma \right\}.$$

Note that C_0 is equal to the real diameter $(-1, 1)$ and for $\gamma \neq 0$, C_γ is a circular arc passing through the points -1 and 1 .

LEMMA 2. *The quantity $|S_{f_0}(z)| \cdot (1 - |z|^2)^2$ is constant on the curves C_γ , $0 \leq \gamma < \pi/2$.*

Proof. Take any $\gamma \in [0, \pi/2)$ and let $z \in C_\gamma$. Then there exists a (positive) real number t such that

$$\frac{1 + z}{1 - z} = te^{i\gamma}.$$

A straightforward computation gives

$$\beta = \frac{t^R e^{iR\gamma}}{t^R \cos(R\gamma)} = 1 + i \tan(R\gamma) \quad \text{and} \quad \frac{1 - |z|^2}{|1 - z^2|} = \cos \gamma.$$

The proof of the lemma follows from (24). \square

Lemma 2 shows, in particular, that whenever $r \in (0, 1)$, $|S_{f_0}(r)|(1 - r^2)^2 = |S_{f_0}(0)| \equiv \lambda$.

Note that $S_{f_0}(\overline{z}) = \overline{S_{f_0}(z)}$ for all $z \in \mathbb{D}$. Moreover, it is easy to check that any radius $\{re^{i\theta}, 0 < r < 1\}$, with $0 < \theta < \pi$, intersects every C_γ with $\gamma > 0$ at exactly one point. Therefore, we conclude that

$$(25) \quad \|S_{f_0}\| = \sup_{0 \leq r < 1} |S_{f_0}(ir)| \cdot (1 - r^2)^2.$$

LEMMA 3. *The Schwarzian norm of f_0 equals*

$$\|S_{f_0}\| = |S_{f_0}(0)| = \lambda.$$

Proof. According to (25), it suffices to show

$$\sup_{0 \leq r < 1} |S_{f_0}(ir)| \cdot (1 - r^2)^2 = |S_{f_0}(0)|.$$

We use (24) to write

$$\begin{aligned} \Phi(r) &= |S_{f_0}(ir)| \cdot (1 - |r|^2)^2 \\ (26) \quad &= \left| 2(1 - a^2) + 2R(R - a)\beta_r - \frac{3R^2\beta_r^2}{2} \right| \cdot \left(\frac{1 - r^2}{1 + r^2} \right)^2 \\ &= |\phi \circ \beta_r| \cdot \left(\frac{1 - r^2}{1 + r^2} \right)^2, \end{aligned}$$

where $\phi(x) = A + Bx + Cx^2$, with $A = 2(1 - a^2)$, $B = 2R(R - a)$, and $C = -3R^2/2$; and, by (23),

$$(27) \quad \beta_r = 1 + i \tan(R\gamma_r) \quad \text{with} \quad \cos \gamma_r = (1 - r^2)/(1 + r^2).$$

We are to check

$$(28) \quad \sup\{\Phi(r) : 0 \leq r < 1\} = \Phi(0).$$

Instead of proving (28), we consider the equivalent problem of showing

$$(29) \quad \sup\{\Phi^2(r) : 0 \leq r < 1\} = \Phi^2(0).$$

The advantage of this new reformulation is that, as the reader may check, we can write

$$|\phi \circ \beta_r|^2 = \tilde{A} + \tilde{B}|\beta_r|^2 + \tilde{C}|\beta_r|^4,$$

with $\tilde{A} = A^2 + 2AB + 4AC$, $\tilde{B} = B^2 + 2BC - 2AC$, and $\tilde{C} = C^2$. Note that $\tilde{A} + \tilde{B} + \tilde{C} = \lambda^2$. In fact,

$$(30) \quad \tilde{A} = 4(1 - a^2)(1 - (a + R)^2)$$

(with $a + R > 1$, by the definition of a),

$$(31) \quad \tilde{B} = 2R^2(3 - aR - a^2 - R^2), \quad \text{and} \quad \tilde{C} = \frac{9R^4}{4}.$$

Notice that by (27), the mapping $r \rightarrow \gamma_r$ is increasing for $r \in (0, 1)$, and that $|\beta_r| = 1/\cos(R\gamma_r)$.

Let us rename $\Psi(r) = \Phi^2(r)$. To prove (29), we distinguish among the following three cases.

(i) $\lambda = 3/2$. In this case, by Theorem 3, $R = 1$. We also have $a = 1$, which gives $A = B = 0$, $C = -3/2$ and therefore, $\tilde{A} = \tilde{B} = 0$, $\tilde{C} = 9/4$. Hence

$$\Psi(r) = \left(\tilde{A} + \tilde{B}|\beta_r|^2 + \tilde{C}|\beta_r|^4 \right) \cdot \cos^4 \gamma_r = \frac{9}{4} \cdot \frac{1}{\cos^4 \gamma_r} \cdot \cos^4 \gamma_r = \frac{9}{4} = \lambda^2.$$

This proves that (29) holds for $\lambda = 3/2$.

(ii) $\lambda > 3/2$. By Theorem 3, we have $R = 1$. A straightforward calculation shows that $a > 1$. Hence, using (30) and that $a + R > 1$, we obtain $\tilde{A} > 0$. Now, since $|\beta_r| = 1/\cos \gamma_r$, we can write

$$\begin{aligned}\Psi(r) &= \left(\tilde{A} + \tilde{B}|\beta|^2 + \tilde{C}|\beta|^4 \right) \cdot \cos^4 \gamma_r = \left(\frac{\tilde{A}}{|\beta|^4} + \frac{\tilde{B}}{|\beta|^2} + \tilde{C} \right) \\ &= \tilde{A} \cos^4(\gamma_r) + \tilde{B} \cos^2(\gamma_r) + \tilde{C}.\end{aligned}$$

Consider the function $\psi(x) = \tilde{A}x^2 + \tilde{B}x + \tilde{C}$, $x \in (0, 1)$. Note that $\Psi(r) = \psi(\cos^2(\gamma_r))$ and that $\Psi(r) \leq \Psi(0)$ if and only if $\psi(x) \leq \psi(1)$. To show that $\psi(x) \leq \psi(1)$ we argue as follows: the graph of ψ is a convex parabola and $\psi(1) = \lambda^2 > 9/4 = \psi(0)$ since $\lambda > 3/2$. Hence, the unique critical point x_0 of ψ is a minimum and satisfies $x_0 < 1$. We have the following possibilities.

- (i) The critical point $x_0 \leq 0$, which implies that for all $r \in [0, 1)$, $\psi(r) \leq \psi(1) = \lambda^2$.
- (ii) The critical point x_0 which is a minimum of ψ belongs to $(0, 1)$. In this case, we have

$$\sup_{0 \leq x \leq 1} \psi(x) = \max\{\psi(0), \psi(1)\} = \max\left\{\frac{9}{4}, \lambda^2\right\} = \lambda^2 = \psi(1),$$

which shows that (29) holds also for $\lambda > 3/2$.

(iii) $\lambda < 3/2$. This is the most complicated case to analyze by far due to the amount of parameters that we are to control. Note that we have that $R < 1$ by Theorem 3, and that $a < 1$ by (19). Hence, using (30) and (31), we have that $\tilde{A} < 0$, while $\tilde{B}, \tilde{C} > 0$. Recall that $\beta_r = 1 + i \tan(R\gamma_r)$ and that $\cos \gamma_r = (1 - r^2)/(1 + r^2)$, which gives that $|\beta_r| = 1/\cos(R\gamma_r)$ and that the correspondence $r \rightarrow \gamma_r$ is an (strictly) increasing function in $r \in (0, 1)$ (and hence $\gamma'_r = \partial \gamma_r / \partial r > 0$ for all such r).

We are to show that $\sup_{0 \leq r < 1} \Psi(r) = \Psi(0)$. To do so, we will check that the derivative of Ψ is non-positive for all such r . For the convenience of the reader, we proceed in different steps.

Step 1: There exists $r_1 \in (0, 1)$ such that $\Psi'(r) < 0$ for all $r \in (0, r_1)$. As it was mentioned before, we can write

$$\Psi(r) = \left(\tilde{A} + \tilde{B}|\beta_r|^2 + \tilde{C}|\beta_r|^4 \right) \cdot \cos^4 \gamma_r.$$

Using this expression for Ψ , we compute its derivative to obtain

$$\begin{aligned}\Psi'(r) &= \left(2\tilde{B}|\beta_r| + 4\tilde{C}|\beta_r|^3 \right) \cdot \frac{\sin(R\gamma_r)}{\cos^2(R\gamma_r)} \cdot R\gamma'_r \cdot \cos^4 \gamma_r \\ &\quad - \left(\tilde{A} + \tilde{B}|\beta_r|^2 + \tilde{C}|\beta_r|^4 \right) \cdot 4 \cos^3 \gamma_r \cdot \sin \gamma_r \cdot \gamma'_r,\end{aligned}$$

which can be written as

$$(32) \quad \Psi'(r) = -\gamma'_r \cdot \frac{\cos^4 \gamma_r \cdot \tan \gamma_r}{\cos(R\gamma_r)} \\ \times \left(4\tilde{A} \cos(R\gamma_r) + 2\tilde{B}|\beta_r| + 2\tilde{B}\varphi_r|\beta_r| + 4\tilde{C}\varphi_r|\beta_r|^3 \right),$$

where

$$\varphi_r = 1 - \frac{R \tan(R\gamma_r)}{\tan \gamma_r}.$$

Let us see that φ_r is an (strictly) increasing function of $r \in (0, 1)$. To do so, we compute the derivative of this function which is

$$\varphi'_r = -\frac{R\gamma'_r}{\tan^2 \gamma_r \cdot \cos \gamma_r \cdot \cos(R\gamma_r)} \cdot \left(\frac{R \sin \gamma_r}{\cos(R\gamma_r)} - \frac{\sin(R\gamma_r)}{\cos \gamma_r} \right).$$

Now, φ_r is increasing if and only if

$$(33) \quad \frac{R \sin \gamma_r}{\cos(R\gamma_r)} < \frac{\sin(R\gamma_r)}{\cos \gamma_r} \quad \Leftrightarrow \quad R \sin(2\gamma_r) - \sin(2R\gamma_r) < 0.$$

Define $\mu(r) = R \sin(2\gamma_r) - \sin(2R\gamma_r)$, $r \in (0, 1)$. Then $\mu(0) = 0$ and

$$\mu'(r) = 2R(\cos(2\gamma_r) - \cos(2R\gamma_r)) \cdot \gamma'_r < 0$$

since $0 < R < 1$, which gives (33) and shows that φ_r is increasing for all $r \in (0, 1)$.

It is also clear that $\cos(R\gamma_r)$ is decreasing in r so that $1/\cos(R\gamma_r) = |\beta_r|$ is increasing as well. Hence, the expression in the parentheses in (32) is an increasing function of $r \in (0, 1)$ (recall that $\tilde{A} < 0$ and that $\tilde{B}, \tilde{C} > 0$). On the other hand, $\cos^4 \gamma_r \cdot \tan \gamma_r$ increases with r for all $0 < r < r_1$, where $\sin(\gamma_{r_1}) = 1/2$. This shows that as long as $0 < r < r_1$, the function

$$\tilde{\varphi}(r) = -\frac{\cos^4 \gamma_r \cdot \tan \gamma_r}{\cos(R\gamma_r)} \cdot \left(4\tilde{A} \cos(R\gamma_r) + 2\tilde{B}|\beta_r| + 2\tilde{B}\varphi_r|\beta_r| + 4\tilde{C}\varphi_r|\beta_r|^3 \right)$$

is decreasing. Since $\tilde{\varphi}(0) = 0$, we get that $\tilde{\varphi}(r) < 0$ for all $r \in (0, r_1)$. In other words, keeping in mind that $\gamma'_r > 0$ for all $r \in (0, 1)$, we conclude from (32) that $\Psi'(r) < 0$ for $r \in (0, r_1)$. (This implies, in particular, that $r = 0$ is a local maximum of Ψ .)

Step 2: There exists $r_2 \in (0, 1)$ such that $\Psi'(r) < 0$ for all $r \in (r_2, 1)$. Recall that

$$\lambda = 2a^2 + 2aR - 2 - R^2/2.$$

Using (26) and (27), we get

$$\begin{aligned}
\Psi(r) &= \left| 2(1 - a^2) + 2R(R - a)\beta_r - \frac{3R^2\beta_r^2}{2} \right|^2 \cdot \left(\frac{1 - r^2}{1 + r^2} \right)^4 \\
&= \left(\operatorname{Re} \left\{ 2(1 - a^2) + 2R(R - a)\beta_r - \frac{3R^2\beta_r^2}{2} \right\} \right)^2 \cdot \cos^4 \gamma_r \\
&\quad + \left(\operatorname{Im} \left\{ 2(1 - a^2) + 2R(R - a)\beta_r - \frac{3R^2\beta_r^2}{2} \right\} \right)^2 \cdot \cos^4 \gamma_r \\
&= \left(-\lambda + \frac{3R^2}{2} \cdot \tan^2(R\gamma_r) \right)^2 \cdot \cos^4 \gamma_r \\
&\quad + (R^2 + 2aR)^2 \cdot \tan^2(R\gamma_r) \cdot \cos^4 \gamma_r \\
(34) \quad &= \lambda^2 \cdot \cos^4 \gamma_r + K \cdot \tan^2(R\gamma_r) \cdot \cos^4 \gamma_r + \frac{9R^4}{4} \cdot \tan^4(R\gamma_r) \cdot \cos^4 \gamma_r,
\end{aligned}$$

where

$$\begin{aligned}
K &= R^4 + 4a^2R^2 + 4aR^3 - 3R^2\lambda = \frac{5R^4}{2} + 6R^2 - 2a^2R^2 - 2aR^3 \\
&\geq \frac{5R^4}{2} + 2R^2 > 0,
\end{aligned}$$

since a and R are less than one. Using (34) we see that the derivative of Ψ with respect to r equals

$$\begin{aligned}
\Psi'(r) &= -4\lambda^2 \cdot \cos^3 \gamma_r \cdot \sin \gamma_r \cdot \gamma'_r + K \cdot \frac{2R \tan(R\gamma_r)}{\cos^2(R\gamma_r)} \cdot \cos^4 \gamma_r \cdot \gamma'_r \\
&\quad - 4K \cdot \cos^3 \gamma_r \cdot \sin \gamma_r \cdot \tan^2(R\gamma_r) \cdot \gamma'_r + 9R^5 \cdot \frac{\tan^3(R\gamma_r)}{\cos^2(R\gamma_r)} \cdot \cos^4 \gamma_r \cdot \gamma'_r \\
&\quad - 9R^4 \tan^4(R\gamma_r) \cdot \cos^3 \gamma_r \cdot \sin \gamma_r \cdot \gamma'_r.
\end{aligned}$$

Note that $\Psi'(r) < 0$ if and only if

$$\begin{aligned}
&K \cdot \frac{2R \tan(R\gamma_r)}{\cos^2(R\gamma_r)} \cdot \cos \gamma_r + 9R^5 \cdot \frac{\tan^3(R\gamma_r)}{\cos^2(R\gamma_r)} \cdot \cos \gamma_r \\
(35) \quad &< 4\lambda^2 \cdot \sin \gamma_r + 4K \cdot \sin \gamma_r \cdot \tan^2(R\gamma_r) + 9R^4 \tan^4(R\gamma_r) \cdot \sin \gamma_r,
\end{aligned}$$

an inequality that holds for all $r \geq r_2$, say, since as $r \rightarrow 1$ (or, equivalently, as $\gamma_r \rightarrow \pi/2$), the right hand side in (35) goes to 0 while the left hand side tends to a strictly positive real number. In other words, this shows that $\Psi'(r) < 0$ for all $r \in (r_2, 1)$.

Step 3: Ψ has at most one critical point in $(0, 1)$. Note that once we check that the number of solutions of the equation $\Psi'(r) = 0$, $r \in (0, 1)$ is at most one, we will have, as a consequence of the previous steps that $\Psi'(r) \leq 0$ for $r \in (0, 1)$ (and hence Ψ is non-increasing in that interval). Observe that $\Psi' = 0$ if and only if the function in the parentheses in (32) is

equal to zero. As was justified in Step 1, this function is increasing, which proves our claim. \square

We summarize the previous analysis in the following proposition.

PROPOSITION 1. *If $\lambda > 0$ then the function $f_0 \in \mathcal{F}_\lambda^0$.*

This shows that (15) and (16) are equalities, and finishes the proof of Theorem 1.

REFERENCES

- [1] M. Chuaqui, P. L. Duren, and B. Osgood, The Schwarzian derivative for harmonic mappings, *J. Anal. Math.* **91** (2003), 329–351.
- [2] M. Chuaqui and B. Osgood, Sharp distortion theorems associated with the Schwarzian derivative, *J. London Math. Soc.* **48** (1993), 289–298.
- [3] J. Clunie and T. Sheil-Small, Harmonic univalent functions, *Ann. Acad. Sci. Fenn. Ser. A* **9** (1984), 3–25.
- [4] P. L. Duren, *Univalent Functions*, Springer-Verlag, New York, 1983.
- [5] P. L. Duren, *Harmonic Mappings in the Plane*, Cambridge University Press, Cambridge, 2004.
- [6] C. Epstein, The hyperbolic Gauss map and quasiconformal reflections, *J. Reine Angew. Math.* **372** (1986), 96–135.
- [7] M. Heins, Some characterizations of finite Blaschke products of positive degree, *J. Anal. Math.* **46** (1986), 162–166.
- [8] R. Hernández and M. J. Martín, Stable geometric properties of analytic and harmonic functions, *Math. Proc. Cambridge Philos. Soc.* **155** (2013), 343–359.
- [9] R. Hernández and M. J. Martín, Quasiconformal extensions of harmonic mappings, *Ann. Acad. Sci. Fenn. Ser. A. I Math.* **38** (2013), 617–630.
- [10] R. Hernández and M. J. Martín, Pre-Schwarzian and Schwarzian derivatives of harmonic mappings, *J. Geom. Anal.* DOI 10.1007/s12220-013-9413-x. Published electronically on April 13th, 2013.
- [11] T. Hosokawa and S. Ohno, Topological structures of the set of composition operators on the Bloch space, *J. Math. Anal. Appl.* **314** (2006), 736–748.
- [12] W. Kraus, Über den Zusammenhang einiger Charakteristiken eines einfach zusammenhängenden Bereiches mit der Kreisabbildung, *Mitt. Math. Sem. Geissen* **21** (1932), 1–28.
- [13] H. Lewy, On the non-vanishing of the Jacobian in certain one-to-one mappings, *Bull. Amer. Math. Soc.* **42** (1936), 689–692.
- [14] M. J. Martín, Hyperbolic distortion of conformal maps at corners, *Constr. Approx.* **30** (2009), 265–275.
- [15] Z. Nehari, The Schwarzian derivative and schlicht functions, *Bull. Amer. Math. Soc.* **55** (1949), 545–551.
- [16] Z. Nehari, A property of convex conformal maps, *J. Anal. Math.* **30** (1976), 390–393.
- [17] Ch. Pommerenke, Linear-invariante Familien analytischer Funktionen I, *Math. Ann.* **155** (1964), 108–154.
- [18] Ch. Pommerenke, Linear-invariante Familien analytischer Funktionen II, *Math. Ann.* **156** (1964), 226–262.
- [19] Ch. Pommerenke, *Univalent Functions*, Vandenhoeck & Ruprecht, Göttingen, 1975.
- [20] T. Sheil-Small, Constants for planar harmonic mappings, *J. London Math. Soc.* **42** (1990), 237–248.

FACULTAD DE MATEMÁTICAS, PONTIFICIA UNIVERSIDAD CATÓLICA DE CHILE, CASILLA 306, SANTIAGO, CHILE.

E-mail address: `mchuaqui@mat.puc.cl`

FACULTAD DE INGENIERÍA Y CIENCIAS, UNIVERSIDAD ADOLFO IBÁÑEZ, AV. PADRE HURTADO 750, VIÑA DEL MAR, CHILE.

E-mail address: `rodrigo.hernandez@uai.cl`

DEPARTMENT OF PHYSICS AND MATHEMATICS, UNIVERSITY OF EASTERN FINLAND, P.O. BOX 111, FI-80101 JOENSUU, FINLAND.

E-mail address: `maria.martin@uef.fi`